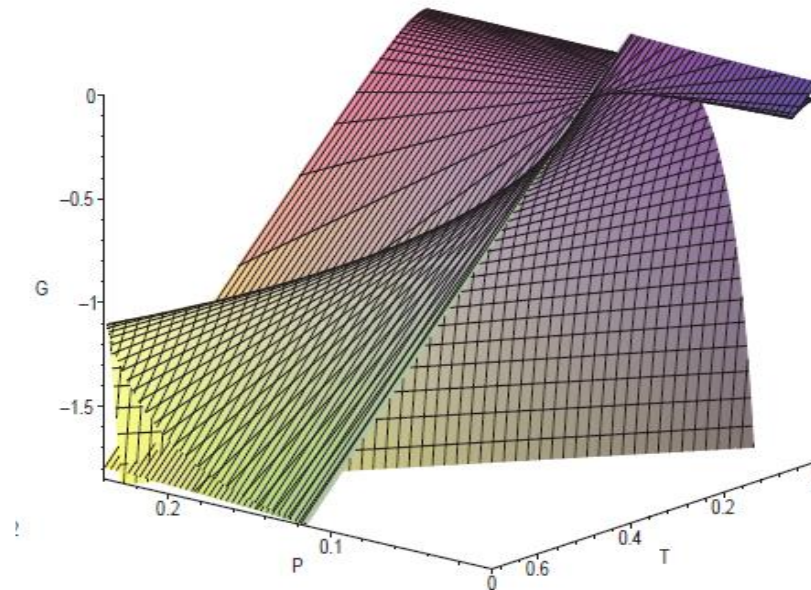


The rise and fall of the black hole chemistry: what happens when we vary Lambda?

David Kubizňák
(ITP, Charles University)



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Gravitation, and Particle Physics**
CEICO, Prague, Czech Republic
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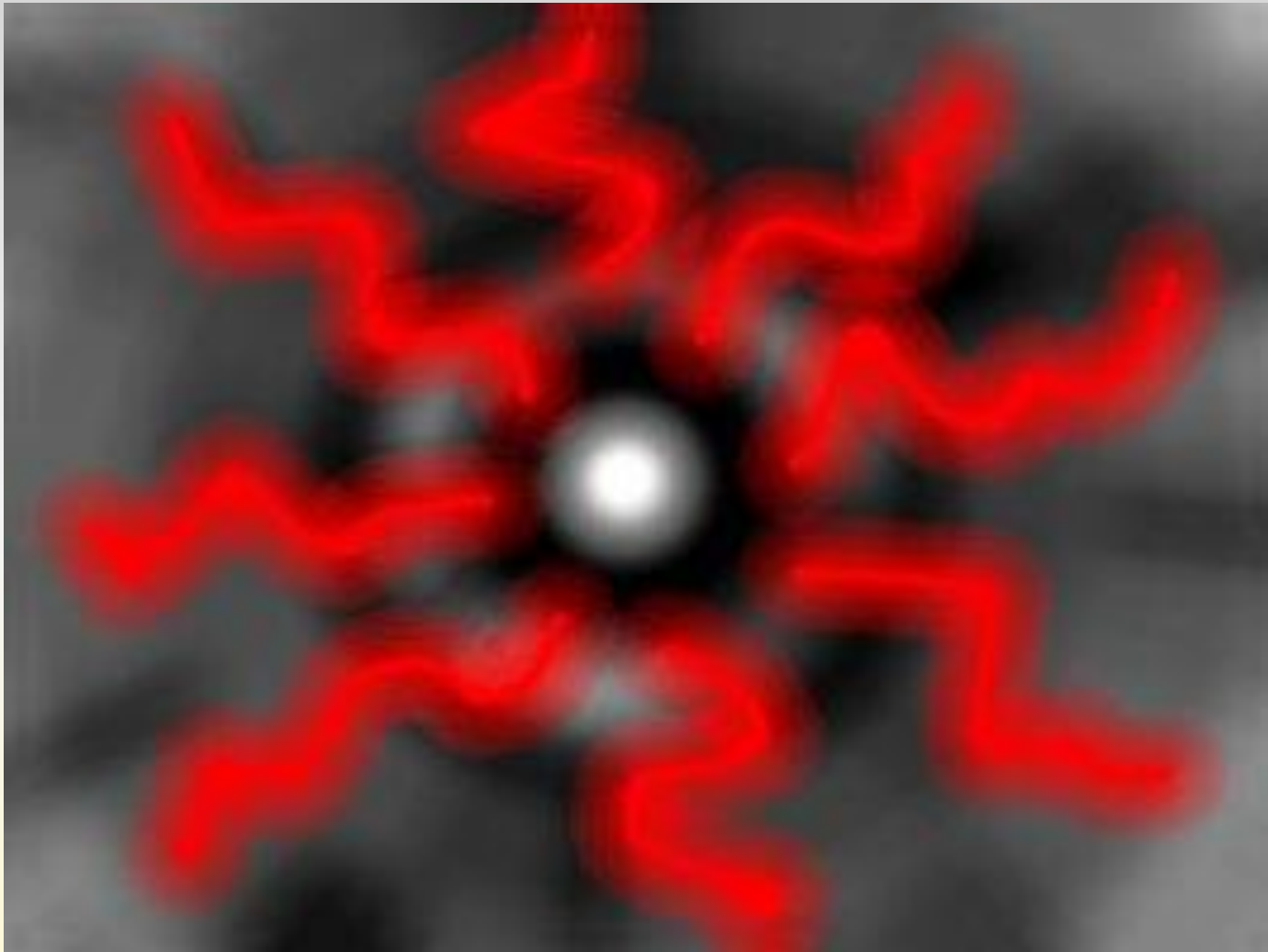
Plan of the talk

- I. Black holes as thermodynamic objects
- II. The rise of black hole chemistry
- III. Extended bulk-boundary correspondence
(the fall of black hole chemistry)
- IV. Summary

Based on:

- DK, R.B. Mann, *P-V criticality of charged AdS black holes*, JHEP 07 (2012) 033; ArXiv:1205:0559.
- W. Cong, DK, R.B. Mann, *Thermodynamics of AdS black holes: central charge criticality*, PRL 127 (2021) 9, 091301; Arxiv:2105.02223.
- W.Cong, DK, R.B. Mann, M.R. Visser, *Holographic CFT phase transitions and criticality of charged AdS black holes*, JHEP 08 (2022) 174.; Arxiv: 2112:14848.

1) Black Holes as Thermodynamic Objects



Black holes as thermodynamic objects

If someone points out to you that your pet theory of the universe is in disagreement with Maxwell's equations-then so much the worse for Maxwell's equations. If it is found to be contradicted by observation-well these experimentalists do bungle things sometimes. But if your theory is found to be against the second law of thermodynamics I can give you no hope; there is nothing for it but to collapse in deepest humiliation.

Sir Arthur Stanley Eddington

Gifford Lectures (1927), *The Nature of the Physical World* (1928), 74.

Black holes and their characteristics

Schwarzschild black hole:

$$ds^2 = -\left(1 - \frac{2M}{r}\right)dt^2 + \frac{dr^2}{1 - \frac{2M}{r}} + r^2 d\Omega^2$$



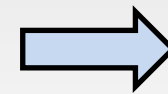
- asymptotic mass (total energy)

$$M = -\frac{1}{8\pi} \int_{S_\infty} *dk, \quad k^a = (\partial_t)^a$$

- black hole horizon: (radius $r_h=2M$)

surface gravity

$$(k^b \nabla_b k^a)|_H = \kappa k^a|_H$$



$$\kappa = \frac{1}{4M}$$

surface area

$$A = 4\pi r_h^2$$

never decreases

$$dM = \kappa dA$$

Bekenstein?

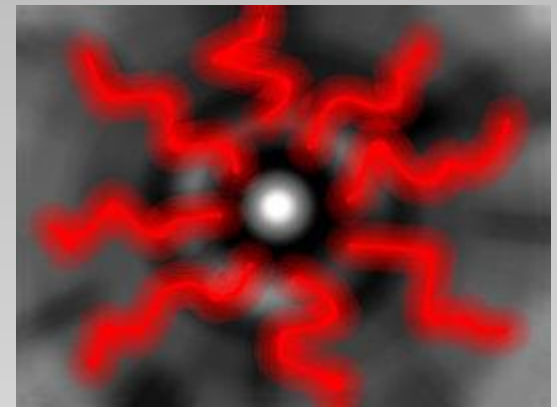


$$dE = T dS$$

Hawking (1974):

$$T = \frac{\kappa}{2\pi}, \quad S = \frac{A}{4}$$

derivation used QFT in curved spacetime

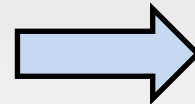


Other approaches:

- Euclidean path integral approach (Gibbons & Hawking-1977)

$$Z = \int D[g_{ab}] e^{-S_E[g]} \approx e^{-S_E[g_c]}$$

$$F = -\frac{1}{\beta} \log Z$$



$$S = -\frac{\partial F}{\partial T} = \frac{A}{4}$$

Euclidean manifold non-singular if the imaginary time τ identified with a certain period $\Delta\tau$. In QFT this corresponds to a finite temperature

$$T = \frac{1}{\beta}, \quad \beta = \Delta\tau$$

- Tunneling approach, LQG, String theory,

Black hole thermodynamics

- First law of black hole thermodynamics:

$$\delta M = T\delta S + \sum_i \Omega_i \delta J_i + \Phi \delta Q$$

- Smarr-Gibbs-Duhem relation:

$$\frac{d-3}{d-2}M = TS + \sum_i \Omega_i J_i + \frac{d-3}{d-2}\Phi Q$$

- Specific heat of AF Schwarzschild BH is negative (cannot have thermal equilibrium)

Where is the PdV term?



2) The rise of
black hole
chemistry

Black hole chemistry

Simple idea:

- Consider an asymptotically **AdS black hole spacetime**
- Identify the cosmological constant with a thermodynamic pressure

$$P = -\frac{\Lambda}{8\pi G}, \quad \Lambda = -\frac{(D-1)(D-2)}{2l^2}$$

- Allow this to be a “dynamical” quantity

(Teitelboim and Brown – 1980’s)

Immediate consequences

- Extended black hole thermodynamics:

D.Kastor, S.Ray, and J.Traschen, *Enthalpy and the Mechanics of AdS Black Holes*, Class. Quant. Grav. 26 (2009) 195011.

$$\delta M = T\delta S + \Theta \delta P + \dots$$

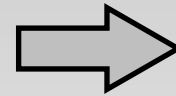
- Introduces the standard **-PdV term** into black hole thermodynamics
- Black hole mass M no longer identified with energy but rather interpreted as **enthalpy**

$$U = M + \epsilon V = M - PV$$

Immediate consequences

- Black hole volume:

$$V = \left(\frac{\partial M}{\partial P} \right)_{S, \dots}$$



Schwarzschild(-AdS):

$$V = \frac{4}{3} \pi r_+^3$$

- More involved for more complicated black holes
- The fact this provides a good definition of volume is supported by the **Reverse Isoperimetric Inequality** conjecture:

M. Cvetič, G.W. Gibbons, D.K. C.N. Pope, *Black hole enthalpy and an entropy inequality for the thermodynamic volume*, Phys. Rev. D84 (2011) 024037, [arXiv:1012.2888].

Immediate consequences

- Consistent Smarr relation:

$$\delta M = T\delta S + V\delta P + \phi\delta Q + \Omega\delta J,$$

$$M = \frac{D-2}{D-3}(TS + \Omega J) + \phi Q - \frac{2}{D-3}PV$$

- Phase transitions:

- AdS black holes can be in **thermal equilibrium**
- Exhibit interesting **phase transitions**
- Provide dual description of CFT at finite temperature via **AdS/CFT correspondence**

Canonical Example: VdW behavior of charged AdS black holes

$$ds^2 = -f dt^2 + \frac{dr^2}{f} + r^2 d\Omega_2^2, \quad A = -\frac{Q}{r} dt$$
$$f = 1 - \frac{2GM}{r} + \frac{GQ^2}{r^2} + \frac{r^2}{l^2},$$

- Basic thermodynamic quantities:

$$M = \frac{r_+(l^2 + r_+^2)}{2l^2G} + \frac{Q^2}{2r_+}, \quad T = \frac{3r_+^4 + l^2r_+^2 - GQ^2l^2}{4\pi l^2 r_+^3}$$
$$S = \frac{\pi r_+^2}{G}, \quad V = \frac{4\pi r_+^3}{3}, \quad \phi = \frac{Q}{r_+},$$

$$F = M - TS = \frac{3GQ^2l^2 + l^2r_+^2 - r_+^4}{4Gr_+l^2}$$

P-V criticality

- DK, R.B. Mann, *P-V criticality of charged AdS black holes*, JHEP 1207 (2012) 033.

Van der Waals fluid

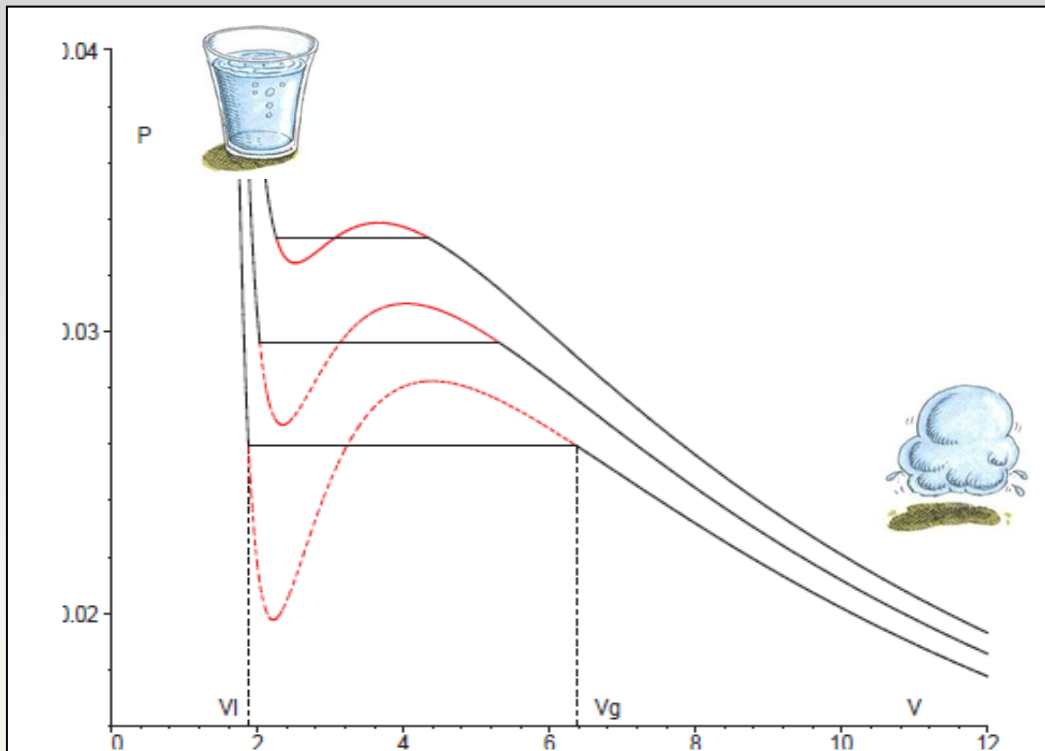


FIG. 2. Maxwell's equal area law. The 'oscillating' (dashed) part of the isotherm $T < T_c$ is replaced by an isobar, such that the areas above and below the isobar are equal one another.

$$\left(P + \frac{a}{v^2}\right) (v - b) = T$$

Parameter a measures the **attraction** between particles ($a > 0$) and b corresponds to "**volume of fluid particles**".

Critical point:

$$\rho_c = \frac{P_c v_c}{T_c} = \frac{3}{8}$$

P-V criticality

- DK, R.B. Mann, *P-V criticality of charged AdS black holes*, JHEP 1207 (2012) 033.

Charged black hole

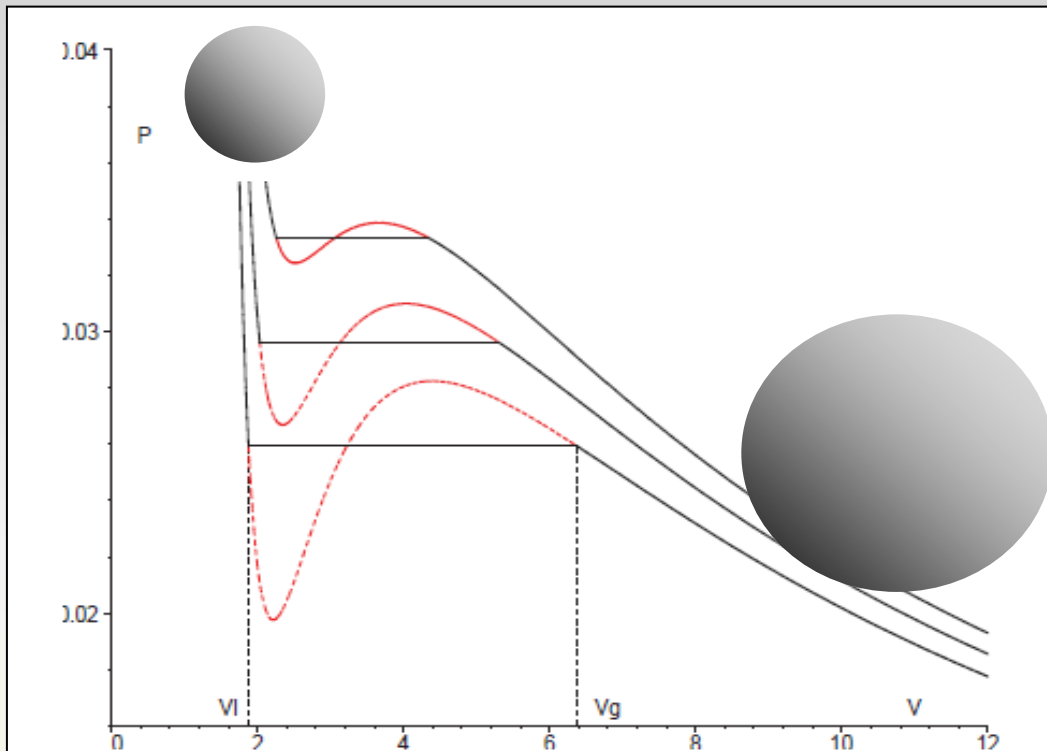


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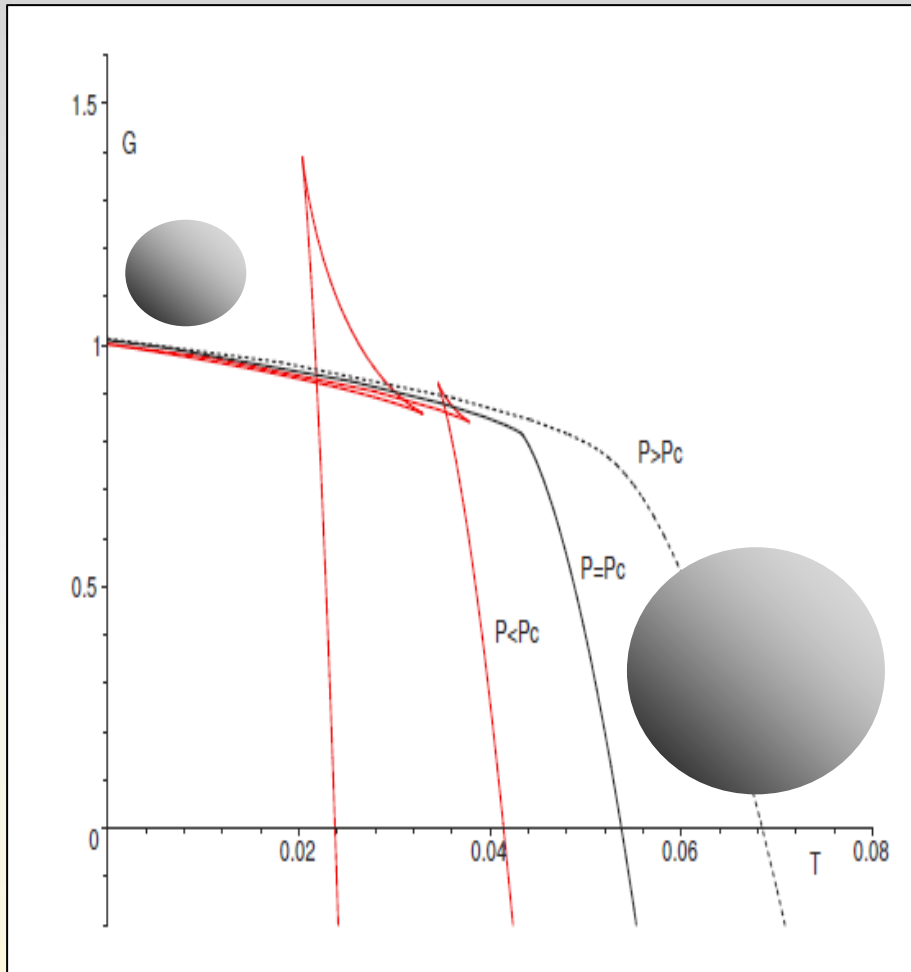
$$P = \frac{T}{v} - \frac{1}{2\pi v^2} + \frac{2Q^2}{\pi v^4}$$

Critical point:

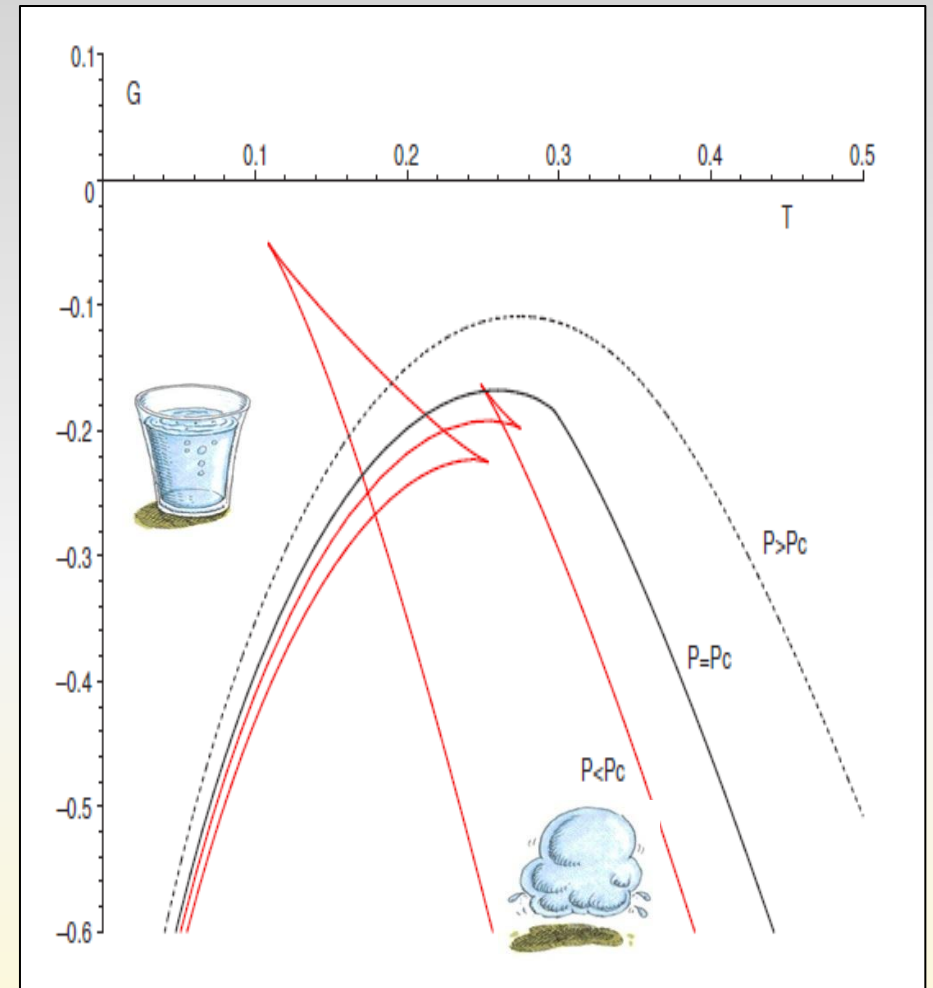
$$\rho_c = \frac{P_c v_c}{T_c} = \frac{3}{8}$$

Free energy: demonstrates standard **swallow tail** behavior

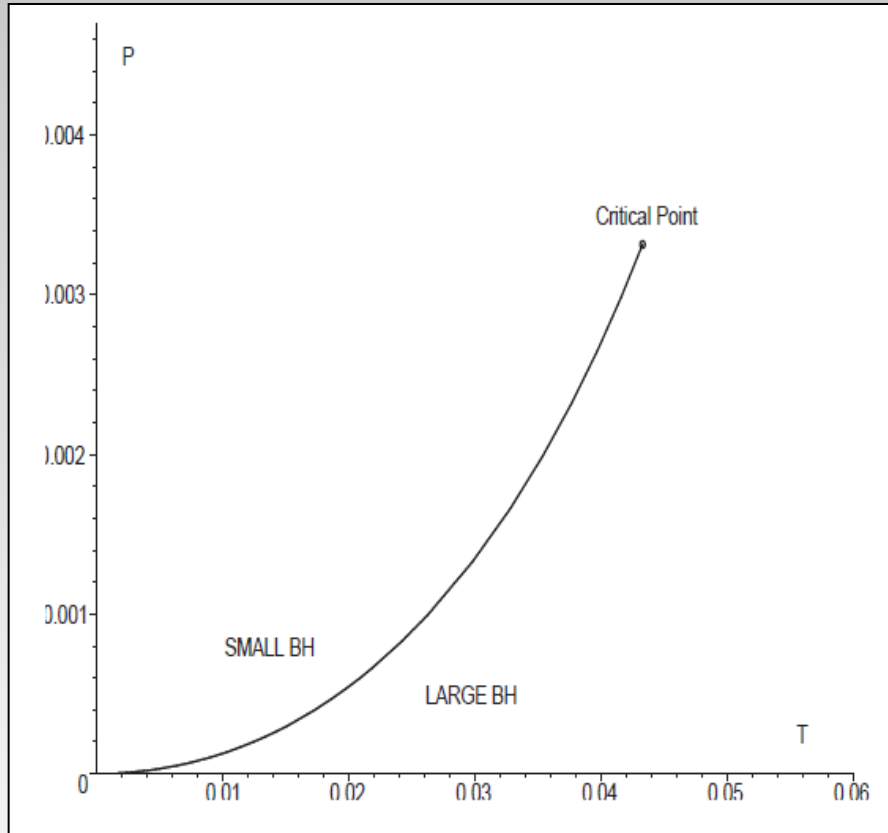
$$F = F(T, P, Q) = \frac{1}{4} \left(r_+ - \frac{8\pi}{3} P r_+^3 + \frac{3Q^2}{r_+} \right)$$



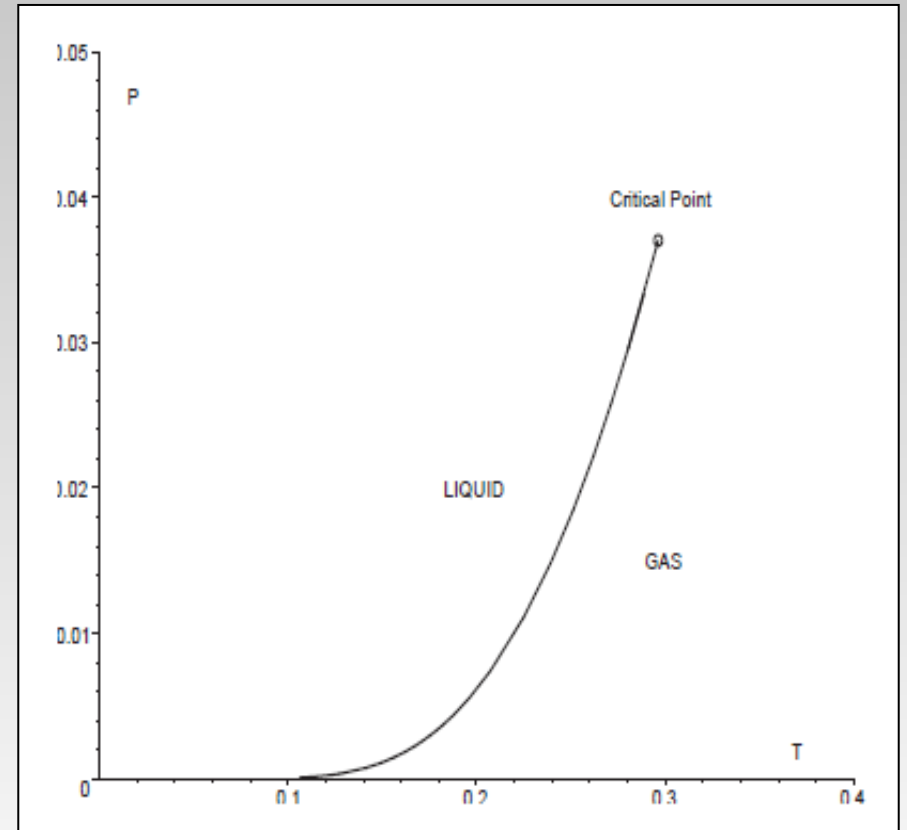
vs.



Phase diagrams: complete analogy



VS.

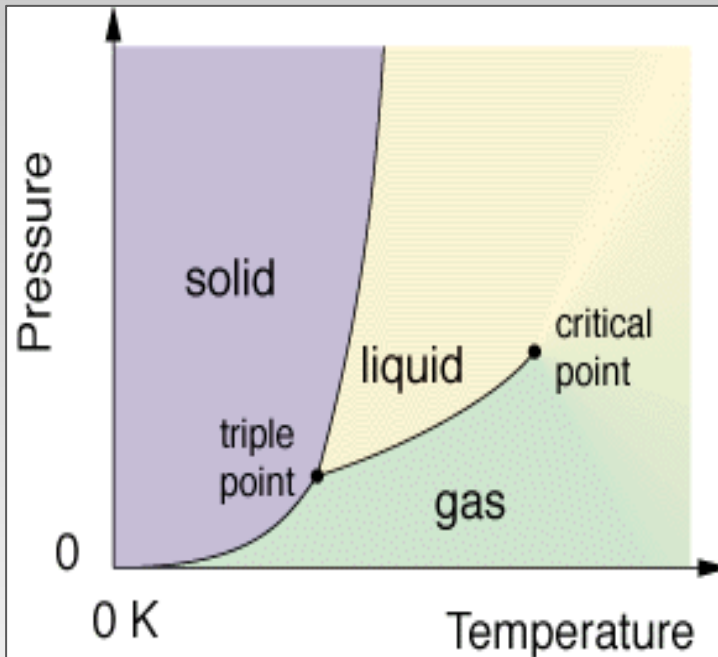


- Coexistence & critical point described by **Clausius-Clapeyron** and **Ehrenfest** equations
- **MFT critical exponents**

$$\alpha = 0, \quad \beta = \frac{1}{2}, \quad \gamma = 1, \quad \delta = 3$$

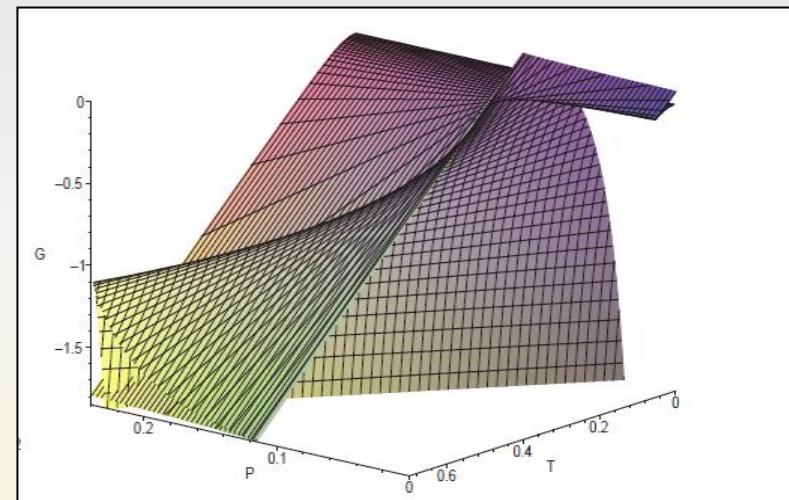
More generally: black hole chemistry

- **Triple point and solid/liquid/gas analogue:**



- **Many other:**

- Isolated critical point
- Reentrant PT
- superfluid PT



- DK, Mann, Teo, *Black hole chemistry: thermodynamics with Lambda*, CQG 34 (2017) 063001, Arxiv:1608.0614.

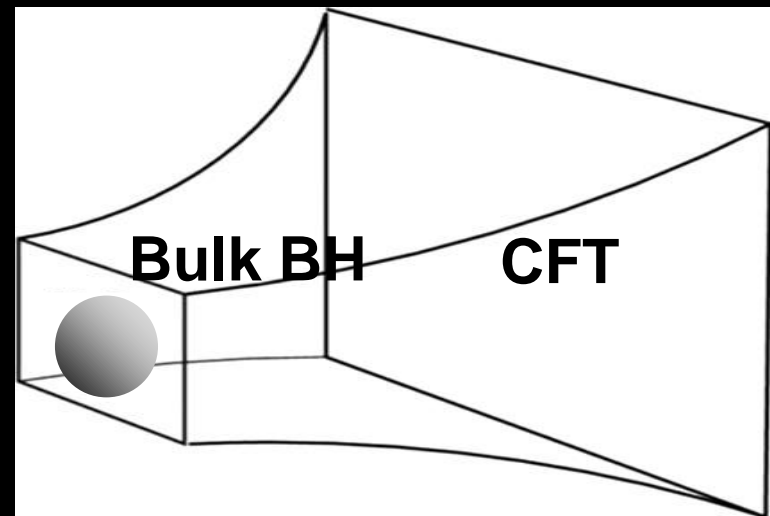
Moral

- We have a very **rich structure** of phase transitions
- **Thermodynamic pressure P** (cosmological constant) plays a role of “**control parameter**”
- But what is the **interpretation on the dual CFT side?**

$$+ V \delta P$$



3) Extended bulk-boundary correspondence (the fall of black hole chemistry)



AdS/CFT interpretation

- Interpretation on boundary field theory

Varying Λ corresponds to varying number of dof N

$$l \sim N$$

e.g. n=4 SU(N) SYM

$$l \propto N^{\frac{1}{4}}$$

$$V \delta P \leftrightarrow \mu \delta N$$

$$C \propto N^2$$

- C. Johnson, Holographic Heat engines, CQG 31 (2014) 205002, Arxiv:1404:5982.
- B. Dolan, Bose condensation and branes, JHEP 10 (2014) 179, ArXiv:1406.7267.
- D.Kastor, S.Ray, J.Traschen, Chemical potential in the first law of holographic entanglement entropy, Arxiv:1409.3521.

AdS/CFT interpretation

- However, this is not the full story, as the **CFT volume** in principle also changes
- A. Karch and B. Robinson, *Holographic black hole chemistry*, JHEP 12 (2015) 073, Arxiv:1510.0247.
- M. Visser, *Holographic thermodynamics requires a chemical potential for color*, Arxiv:2101.04145.

$$ds^2 = -\frac{r^2}{L^2} dt^2 + \frac{L^2}{r^2} dr^2 + r^2 d\Omega_{d-1}^2$$

$$\lambda = R/r$$

$$ds^2 = -\frac{R^2}{L^2} dt^2 + R^2 d\Omega_{d-1}^2$$

AdS/CFT interpretation

- **Standardly**, one chooses

$$R = L \quad \Rightarrow \quad \mathcal{V} = \Omega_{d-1} L^{d-1}$$

and the volume also changes with variations of L .

- **The second** (quite natural from CFT side) **possibility** is to keep R general.

In this case, variations of L are independent of variations of the CFT volume. However, what do the variations of R correspond to in the bulk?

- In any case, we have the following CFT TDs:

Holographic first law (Visser 21)

$$\delta E = T\delta S - pd\mathcal{V} + \tilde{\phi}\delta\tilde{Q} + \Omega\delta J + \tilde{\mu}\delta C$$

- Here, E is the **internal energy** not **enthalpy**!
- Allows to study μ -C criticality

- Accompanied by $E = (D - 2)p\mathcal{V}$

$$E = TS + \tilde{\phi}\tilde{Q} + \Omega J + \tilde{\mu}C$$

- The first equality - **equation of state** - comes from the dimensional analysis
- The second equality – **holographic Smarr** – comes from “central charge extensivity” (note that it does not have D-dependent factors!)

Corresponding bulk first law

- Holographic dictionary

$$C = \frac{\Omega_{d-1} L^{d-1}}{16\pi G_N} \quad (\text{Einstein gravity})$$

$$\tilde{Q} = QL.$$

$$S = \frac{A}{4G_N},$$

$$E = M \frac{L}{R},$$

$$T = \frac{\kappa}{2\pi} \frac{L}{R}$$

$$\tilde{\Phi} = \frac{\Phi L}{L R},$$

- Yields the following **bulk first law**:

$$dM = \frac{\kappa}{8\pi G_N} dA + \Phi dQ + \frac{\Theta}{8\pi G_N} d\Lambda - (M - \Phi Q) \frac{dG_N}{G_N}$$

- Reduces to the previous when G fixed
- Note also, that one of dG or dΛ is not “required” if R also varied.

Bulk first law with “mixed variables”

- Starting with the above bulk first law:

$$\delta M = \frac{\kappa}{8\pi G} \delta A + \Omega \delta J + \phi \delta Q - \frac{V}{8\pi G} \delta \Lambda - \alpha \frac{\delta G}{G}$$

- We can change variables and obtain a Mixed first law:

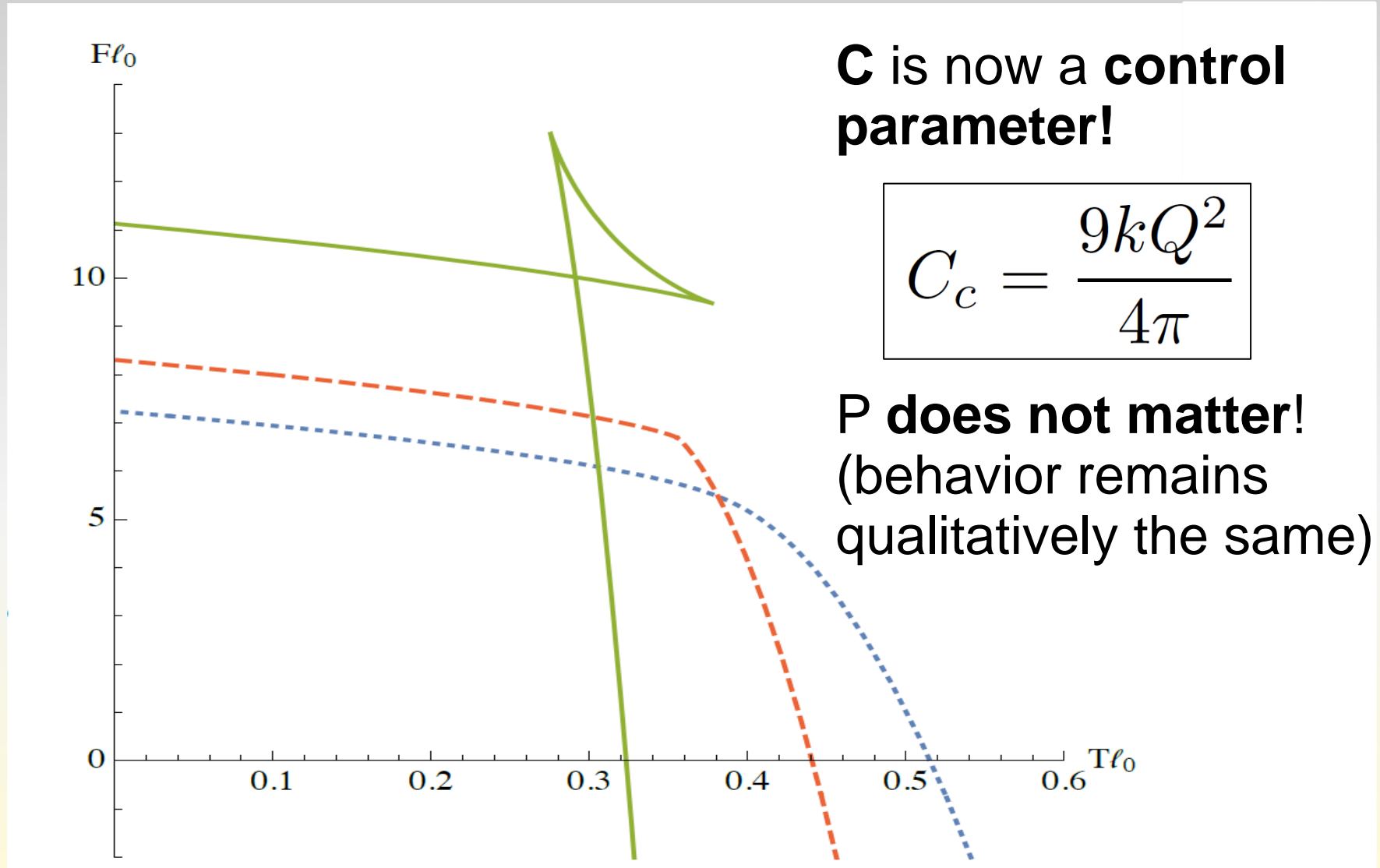
$$C = k \frac{l^{D-2}}{16\pi G} \quad \Rightarrow \quad \frac{\delta G}{G} = -\frac{2}{D} \frac{\delta C}{C} - \frac{D-2}{D} \frac{\delta P}{P}$$

$$\delta M = T \delta S + \Omega \delta J + \phi \delta Q + V_C \delta P + \mu \delta C$$

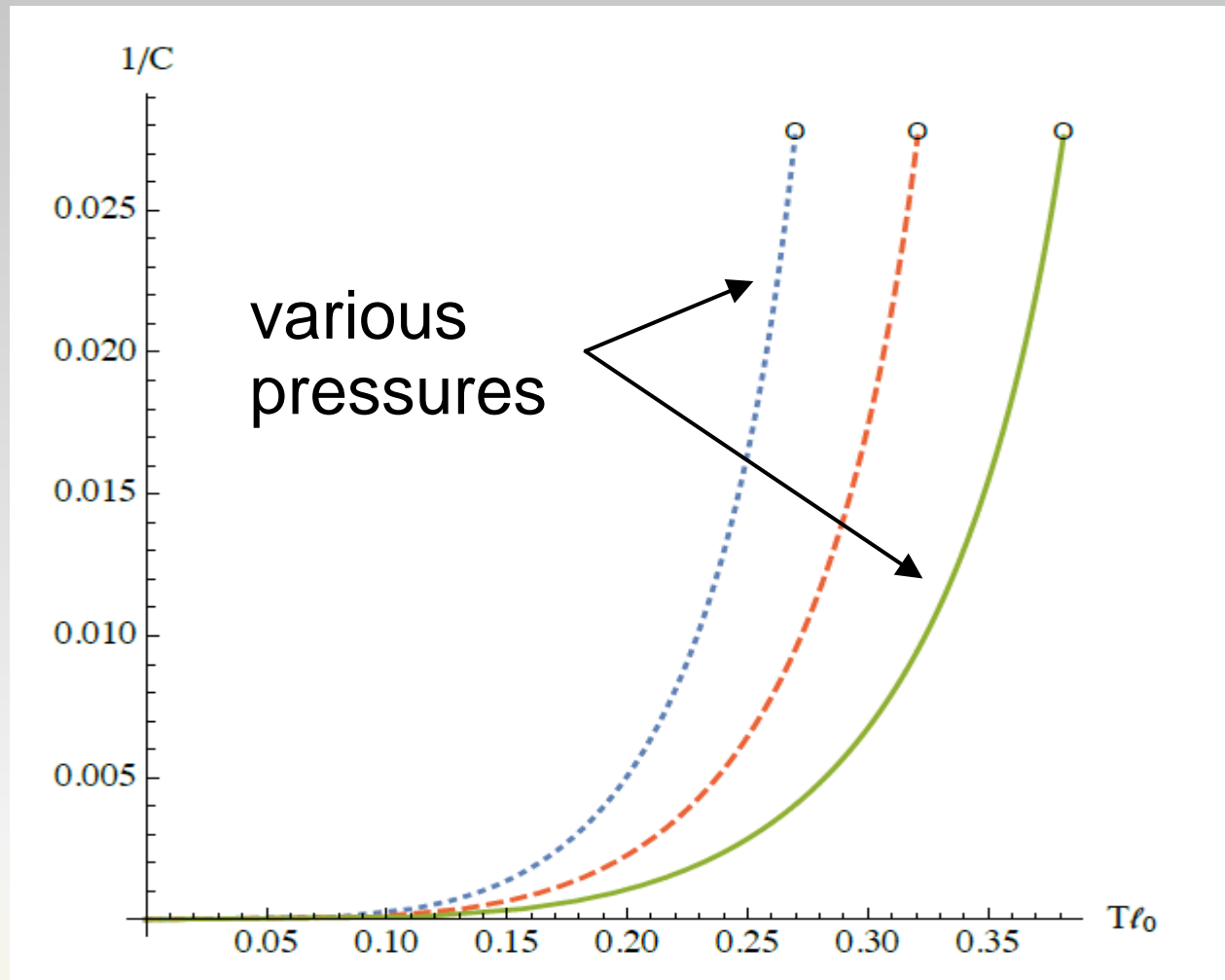
- Enables one to study **bulk** TDs at **fixed central charge C** (while both G and l vary)

Bulk mu-C criticality

$$F = M - TS = F(T, P, Q, C)$$



Phase diagram



Only when CFT has **large number of dof**, $C > C_c$, the bulk black hole can experience a phase transition. Thus, **BH chemistry interpretation seems lost!**

Summary

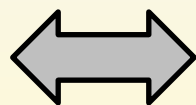
- 1) **Black hole chemistry** (TDs with variable Lambda) provides an interesting framework for AdS black hole thermodynamics:

$$\delta M = T\delta S + V\delta P + \phi\delta Q + \Omega\delta J$$

- **Extended first law** and consistent Smarr
- Black hole mass is **enthalpy**
- Definition for **black hole volume**
- Uncovers various **phase transitions** & similarities with TDs of ordinary systems:

e.g. **Van der Waals criticality** of charged BHs

- 2) On CFT side, **varying Lambda** corresponds to both – varying the **central charge C** and varying the **CFT volume**.

 $\delta\Lambda$  δC

&

 $\delta\mathcal{V}$

Summary

3) The full **CFT first law**:

$$\delta E = T\delta S - pd\mathcal{V} + \tilde{\phi}\delta\tilde{Q} + \Omega\delta J + \tilde{\mu}\delta C$$

standardly translates to **varying both** G and Λ in the bulk!

- However, by considering CFT on a sphere of **radius R** , it is possible to vary only one of these (e.g. Λ)
- This then corresponds to the “**standard black hole chemistry**” with its nice bulk interpretation, and its “**old**” **holographic interpretation** with variable central charge.

Summary

- 4) For the bulk thermodynamics (with variable G and Λ) we can write the following **mixed law**:

$$\delta M = T\delta S + \Omega\delta J + \phi\delta Q + V_C\delta P + \mu\delta C$$

This allows to study TDs of bulk black holes for fixed **CFT central charge C** (kind of TD ensemble).

Interestingly, thermodynamic pressure P is no longer a “control parameter” – **criticality only depends** on the **number of dof** of the dual CFT.

The nice interpretation of **black hole chemistry** **seems to fall** in this case!