Symmetric Teleparallel Horndeski

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Based on: 2212.08005

16/12/2022

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2022 Winter CAS-JSPS Workshop in Cosmology, Gravitation, and Particle Physics, Prague

Outline

- Symmetric teleparallel gravity
- Symmetric teleparallel Horndeski: scheme
- Systematic construction
- Cosmology
- Conclusion

Geometric preliminaries

J.D. McCrea 92.

• Fundamental object: non-metricity tensor $Q_{\lambda\mu\nu} \equiv \nabla_{\lambda}g_{\mu\nu} = \partial_{\lambda}g_{\mu\nu} - \Gamma^{\rho}{}_{\lambda\mu}g_{\rho\nu} - \Gamma^{\rho}{}_{\lambda\nu}g_{\mu\rho}$



Symmetric teleparallel (ST) gravity

• No torsion: $T^{\lambda}{}_{\mu\nu} \equiv \Gamma^{\lambda}{}_{\mu\nu} - \Gamma^{\lambda}{}_{\nu\mu} = 0$, $\Gamma^{\alpha}{}_{\mu\nu} = (\Lambda^{-1})^{\alpha}{}_{\lambda}\partial_{\mu}\Lambda^{\lambda}{}_{\nu}$

• No curvature:
$$R^{\alpha}_{\ \rho\mu\nu} \equiv \partial_{\mu}\Gamma^{\alpha}_{\ \nu\rho} - \partial_{\nu}\Gamma^{\alpha}_{\ \mu\rho} + \Gamma^{\alpha}_{\ \mu\beta}\Gamma^{\beta}_{\ \nu\rho} - \Gamma^{\alpha}_{\ \nu\beta}\Gamma^{\beta}_{\ \mu\rho} = 0 \qquad \Longrightarrow \qquad \Gamma^{\alpha}_{\ \mu\nu} = \frac{\partial x^{\alpha}}{\partial \xi^{\lambda}}\partial_{\mu}\partial_{\nu}\xi^{\lambda}$$

• Can relate
$$R^{\sigma}_{\rho\mu\nu} = \mathring{R}^{\sigma}_{\rho\mu\nu} - \mathring{\nabla}_{\nu}L^{\sigma}_{\mu\rho} + \mathring{\nabla}_{\mu}L^{\sigma}_{\nu\rho} - L^{\sigma}_{\nu\lambda}L^{\lambda}_{\mu\rho} + L^{\sigma}_{\mu\lambda}L^{\lambda}_{\nu\rho} = 0$$

- Coincident gauge $\xi^{\mu} = x^{\mu} \Rightarrow \Gamma^{\alpha}{}_{\mu\nu} = 0 \Rightarrow \nabla_{\mu} = \partial_{\mu}$
- Einstein-Hilbert action becomes "Einstein action"; diff-invariance lost.

$$S_{\rm STEGR}|_{\Gamma=0} = \frac{M_{\rm pl}^2}{2} \int d^4x \sqrt{-g} g^{\mu\nu} \left(\mathring{\Gamma}^{\alpha}{}_{\beta\mu} \mathring{\Gamma}^{\beta}{}_{\nu\alpha} - \mathring{\Gamma}^{\alpha}{}_{\beta\alpha} \mathring{\Gamma}^{\beta}{}_{\mu\nu} \right)$$

2 formulations of GR

J. B. Jiménez et al. '19

• Riemannian GR: Curvature
$$R^{\alpha}_{\ \rho\mu\nu} = \mathring{R}^{\alpha}_{\ \rho\mu\nu} \neq 0$$
 $T^{\lambda}_{\ \mu\nu} = 0$ $Q_{\lambda\mu\nu} = 0$

$$S = \frac{M_{\rm Pl}^2}{2} \int d^4x \sqrt{-g} \mathring{R} \qquad \Longrightarrow \qquad 10 - 2x4 = 2 \text{ propagating d.o.f}$$

• STEGR: Non-metricity

Equivalently
$$Q_1 = W^{\mu}W_{\mu}$$

could be
expressed in $Q_2 = \Lambda_{\mu}\Lambda^{\mu}$
terms of $Q_3 = W_{\mu}\Lambda^{\mu}$
 $Q_4 = *\Omega_{\alpha\mu\nu}*\Omega^{\alpha\mu\nu}$
 $Q_5 = q_{\lambda\mu\nu}q^{\lambda\mu\nu}$

$$\begin{split} S_{\text{Newer GR}} &= \frac{M_{\text{Pl}}^2}{2} \int d^4 x \sqrt{-g} \, \mathbb{Q} \\ \mathbb{Q} &= c_1 Q_{\alpha}^{\ \mu\nu} Q_{\mu\nu}^{\alpha} + c_2 Q_{\alpha}^{\ \mu\nu} Q_{\mu\nu}^{\ \alpha} + c_3 Q_{\mu} Q^{\mu} + c_4 \tilde{Q}_{\mu} \tilde{Q}^{\mu} + c_5 \tilde{Q}_{\mu} Q^{\mu} \\ c_1 &= \frac{1}{4} \,, \quad c_2 &= -\frac{1}{2} \,, \quad c_3 &= -\frac{1}{4} \,, \quad c_4 &= 0 \,, \quad c_5 &= \frac{1}{2} \quad \Longrightarrow \quad \text{Non-metricity scalar} \\ \mathring{R} &= -Q + \mathring{\nabla}_{\mu} (\tilde{Q}^{\mu} - Q^{\mu}) := -Q + B_Q \quad \Longrightarrow \quad \text{Equivalent up to boundary terms} \end{split}$$

ST Horndeski - conditions

- Euler-Lagrange (EL) equations for all the fields $(g_{\mu\nu}, \Gamma^{\lambda}{}_{\mu\nu}, \phi)$ at most second order
 - ➡ No Ostrogradsky ghosts
- Lagrangian is parity preserving \implies useful to define $*\Omega^{\alpha\mu\nu} = \epsilon^{\rho\sigma\mu\nu}\Omega^{\alpha}{}_{\rho\sigma}$
- Invariants at most quadric in the non-metricity $Q_{\lambda\mu\nu}$
- Max. # of d.o.f. = 11 (easy to see in coincident gauge)

Covariantization & variations

• Covariantization prescription: $\eta_{\mu\nu} \rightarrow g_{\mu\nu}$

$$\partial_{\mu} \rightarrow \mathring{\nabla}_{\mu}$$

- In ST-geometry: in principle other prescription possible: $\eta_{\mu\nu} \rightarrow g_{\mu\nu}$ $\partial_{\mu} \rightarrow \nabla_{\mu}$
 - Matter couples non-minimally
- Variation w.r.t. the metric, the scalar field and the flat connection
- 2 ways of ST extensions:
 - a) Dependence of G_i functions on $Q_{\lambda\mu\nu}$ (1st order derivatives)
 - b) Derivative terms $\mathring{
 abla}_{\alpha}Q_{\lambda\mu\nu}$ in Lagrangian (2nd order derivatives)



- 2nd order EL eqs. not guaranteed yet
- Full Lagrangian: $L = \sum_{i=2} c_i L_i$

ST Horndeski – special cases

• Q only:
$$L = \sum_{N_Q \ge 0} C_{N_Q} \mathcal{Q}^{N_Q} \equiv f(\{\mathcal{Q}\})$$

 $N_{\varphi} = 0$



$$L = f(Q, Q_1, Q_2, Q_3, Q_4)$$

i = 2

• Riemannian Horndeski:
$$\mathring{L}_{i} = \sum_{N_{\phi} \ge n \ge 0} C_{N_{\phi},n} \left[\sum_{\substack{m,r \ge 0 \\ m+2r=i-2}} A_{m,r}^{(n,N_{\phi})} \phi^{N_{\phi}-n-m} (\partial \phi)^{n} (\mathring{\nabla} \mathring{\nabla} \phi)^{m} \mathring{\mathcal{R}}^{r} \right]$$

 $2 \leq i \leq 5$

L₂– systematic construction (no HD)

• General form:
$$L_{2} = \sum_{N_{\phi}, N_{Q} \ge 0} \sum_{n=0}^{N_{\phi}} c_{N_{\phi}, N_{Q}, n} \phi^{N_{\phi} - n} (\partial \phi)^{n} \mathcal{Q}^{N_{Q}} \equiv L_{2}(\phi, \partial \phi, \{Q\}) \qquad \mathcal{Q} = (W^{\mu}, \Lambda^{\mu}, q^{\mu\nu\rho}, *\Omega^{\alpha\mu\nu})$$
$$\boxed{\mathbf{n} = \mathbf{1} \quad I_{1} = W^{\mu}\phi_{;\mu}, \\ I_{2} = \Lambda^{\mu}\phi_{;\mu}, \\ J_{2} = q_{\lambda\mu\nu}\Lambda^{\lambda}\phi^{;\mu}\phi^{;\nu}, \\ J_{3} = *\Omega_{\mu}^{\nu\sigma}M_{\sigma}\phi^{;\mu}\phi_{;\nu}, \\ J_{4} = *\Omega_{\mu}^{\nu\sigma}\Lambda_{\sigma}\phi^{;\mu}\phi^{;\nu}\phi_{;\nu}, \\ J_{5} = q_{\lambda\alpha\mu}q^{\lambda\alpha}{}_{\nu}\phi^{;\mu}\phi^{;\nu}, \\ J_{6} = q_{\lambda\alpha\mu}*\Omega^{\lambda\alpha}{}_{\nu}\phi^{;\mu}\phi^{;\nu}, \\ J_{7} = *\Omega_{\lambda\alpha\mu}*\Omega^{\lambda\alpha}{}_{\nu}\phi^{;\mu}\phi^{;\nu}, \\ J_{8} = *\Omega_{\lambda\alpha\mu}*\Omega_{\nu}^{\lambda\alpha}\phi^{;\mu}\phi^{;\nu}. \end{aligned}$$
$$\boxed{\mathbf{n} = \mathbf{1} \quad \mathbf{n} \quad \mathbf{n} = \mathbf{1} \quad \mathbf{n} \quad \mathbf$$

 $L_2 = \mathring{L}_2 + L_{2-\text{STele}} = G_2(\phi, X) + G_{\text{ST}}(Q, Q_1, Q_2, Q_3, Q_4, \phi, X, I_1, I_2, J_1, J_2, J_3, J_4, J_5, J_6, J_7, J_8, J_9, J_{10}, J_{11}, J_{12})$

L₃– systematic construction (HD)

$$\text{General form:} \quad L_3 = \sum_{\substack{N_{\phi}, N_Q \ge 0 \\ N_{\phi} \ge n \ge 0}} C_{N_{\phi}, N_Q, n} \left[\sum_{\substack{m,l \ge 0 \\ m+l=1}} A_{m,l}^{(n,N_{\phi},N_Q)} \ \phi^{N_{\phi}-n-m} (\partial \phi)^n \ (\mathring{\nabla} \mathring{\nabla} \phi)^m \ \mathcal{Q}^{N_Q-l} \ (\mathring{\nabla} \mathcal{Q})^l \right]$$

$$\begin{split} \mathbf{N}_{\mathbf{Q}} &= \mathbf{0} \\ & \mathring{L}_{3} = L_{3}(N_{Q} = 0) = \sum_{\substack{N_{\phi} \ge 0 \\ N_{\phi} \ge n \ge 0}} C_{N_{\phi}, N_{Q} = 0, n} \left[A_{1,0}^{(n,N_{\phi},N_{Q}=0)} \phi^{N_{\phi}-n-1} (\partial \phi)^{n} (\mathring{\nabla} \mathring{\nabla} \phi) \right] \\ &= -G_{3}(\phi, X) \mathring{\Box} \phi, \end{split}$$
 Riemannian KGB

$$\begin{split} \mathbf{N}_{\mathbf{Q}} &= \mathbf{1} \\ L_{3}(N_{Q} = 1) = \sum_{\substack{N_{\phi} \ge 0 \\ N_{\phi} \ge n \ge 0}} C_{N_{\phi}, N_{Q} = 1, n} \left[A_{1,0}^{(n, N_{\phi}, N_{Q} = 1)} \phi^{N_{\phi} - n - 1} (\partial \phi)^{n} (\mathring{\nabla} \mathring{\nabla} \phi) \mathcal{Q} + A_{0,1}^{(n, N_{\phi}, N_{Q} = 1)} \phi^{N_{\phi} - n} (\partial \phi)^{n} (\mathring{\nabla} \mathcal{Q}) \right] \\ &= \sum_{a} \left[\tilde{G}_{3}^{(a)}(\phi, X) \tilde{\mathcal{O}}_{a} + F_{3}^{(a)}(\phi, X) \hat{\mathcal{O}}_{a} \right], \quad \Longrightarrow \quad 2 \text{ types of contractions; not independent.} \end{split}$$

L₃ – systematic construction (HD)

• All possible contractions linear in Q and $\mathring{\nabla}\mathring{\nabla}\phi$ with increasing factors of $\mathring{\nabla}\phi$:

$$\begin{split} \tilde{\mathcal{O}}_{1} &= \qquad \phi^{\mu} Q_{\mu\nu} {}^{\nu} \overset{\circ}{\Box} \phi \qquad = 4\mathcal{O}_{W1} \,, \\ \tilde{\mathcal{O}}_{2} &= \qquad \phi^{\mu} Q_{\nu\mu} {}^{\nu} \overset{\circ}{\Box} \phi \qquad = \frac{9}{4} \mathcal{O}_{\Lambda 1} + \mathcal{O}_{W1} \,, \\ \tilde{\mathcal{O}}_{3} &= \qquad \phi_{;\alpha} Q_{\beta\mu} {}^{\mu} \overset{\circ}{\nabla} {}^{\alpha} \overset{\circ}{\nabla} {}^{\beta} \phi \qquad = 4\mathcal{O}_{W2} \,, \\ \tilde{\mathcal{O}}_{4} &= \qquad \phi_{;\alpha} Q_{\mu\beta} {}^{\mu} \overset{\circ}{\nabla} {}^{\alpha} \overset{\circ}{\nabla} {}^{\beta} \phi \qquad = \mathcal{O}_{W2} + \frac{1}{4} \mathcal{O}_{\Lambda 2} + \frac{1}{2} \mathcal{O}_{\Lambda 1} + \frac{1}{6} \mathcal{O}_{\Omega 1} + \mathcal{O}_{q1} \,, \\ \tilde{\mathcal{O}}_{5} &= \qquad \phi^{;\mu} Q_{\alpha\mu\beta} \overset{\circ}{\nabla} {}^{\alpha} \overset{\circ}{\nabla} {}^{\beta} \phi \qquad = \mathcal{O}_{W2} + \frac{1}{4} \mathcal{O}_{\Lambda 2} - \frac{1}{4} \mathcal{O}_{\Lambda 1} - \frac{1}{3} \mathcal{O}_{\Omega 1} + \mathcal{O}_{q1} \,, \\ \tilde{\mathcal{O}}_{6} &= \qquad \phi^{;\mu} Q_{\mu\alpha\beta} \overset{\circ}{\nabla} {}^{\alpha} \overset{\circ}{\nabla} {}^{\beta} \phi \qquad = \mathcal{O}_{W1} + \mathcal{O}_{\Lambda 2} - \frac{1}{4} \mathcal{O}_{\Lambda 1} - \frac{1}{3} \mathcal{O}_{\Omega 1} + \mathcal{O}_{q1} \,, \\ \tilde{\mathcal{O}}_{7} &= \qquad \phi^{;\mu} \phi^{;\nu} \phi^{;\alpha} Q_{\mu\nu\alpha} \overset{\circ}{\Box} \phi \qquad = 4\mathcal{O}_{W3} \,, \\ \tilde{\mathcal{O}}_{8} &= \qquad \phi^{;\mu} \phi^{;\nu} \phi^{;\alpha} Q_{\mu\beta} {}^{\beta} \overset{\circ}{\nabla}_{\mu} \overset{\circ}{\nabla}_{\alpha} \phi \qquad = 4\mathcal{O}_{W3} \,, \\ \tilde{\mathcal{O}}_{9} &= \qquad \phi^{;\mu} \phi^{;\nu} \phi^{;\alpha} Q_{\mu\nu\beta} {}^{\beta} \overset{\circ}{\nabla}_{\mu} \overset{\circ}{\nabla}_{\alpha} \phi \qquad = \frac{9}{4} \mathcal{O}_{\Lambda 3} + \mathcal{O}_{W3} \,, \\ \tilde{\mathcal{O}}_{10} &= \qquad \phi^{;\mu} \phi^{;\nu} \phi^{;\alpha} Q_{\mu\nu\beta} {}^{\beta} \overset{\circ}{\nabla}_{\alpha} \overset{\circ}{\nabla}_{\beta} \phi \qquad = -2X \, \mathcal{O}_{W2} + \mathcal{O}_{\Lambda 3} + \frac{1}{6} \mathcal{O}_{\Omega 2} + \mathcal{O}_{q3} \,, \\ \tilde{\mathcal{O}}_{11} &= \qquad \phi^{;\mu} \phi^{;\nu} \phi^{;\alpha} Q^{\beta} {}^{\mu} {}^{\nu} \overset{\circ}{\nabla}_{\alpha} \overset{\circ}{\nabla}_{\beta} \phi \qquad = -2X \, \mathcal{O}_{W3} - \frac{3}{2} X \, \mathcal{O}_{\Lambda 3} + \mathcal{O}_{q4} \,, \end{aligned}$$

$$\begin{aligned} \mathcal{O}_{W1} &= W_{\mu}\phi^{;\mu} \mathring{\Box}\phi, \\ \mathcal{O}_{W2} &= W_{\alpha}\phi_{;\beta} \mathring{\nabla}^{\alpha} \mathring{\nabla}^{\beta}\phi, \\ \mathcal{O}_{W3} &= W_{\mu}\phi^{;\mu}\phi^{;\alpha}\phi^{;\beta} \mathring{\nabla}_{\alpha} \mathring{\nabla}_{\beta}\phi, \\ \mathcal{O}_{\Lambda1} &= \Lambda_{\mu}\phi^{;\mu} \mathring{\Box}\phi, \\ \mathcal{O}_{\Lambda2} &= \Lambda_{\alpha}\phi_{;\beta} \mathring{\nabla}^{\alpha} \mathring{\nabla}^{\beta}\phi, \\ \mathcal{O}_{\Lambda3} &= \Lambda_{\mu}\phi^{;\mu}\phi^{;\alpha}\phi^{;\beta} \mathring{\nabla}_{\alpha} \mathring{\nabla}_{\beta}\phi, \\ \mathcal{O}_{\Omega1} &= *\Omega_{\alpha\beta\mu}\phi^{;\mu} \mathring{\nabla}^{\alpha} \mathring{\nabla}^{\beta}\phi, \\ \mathcal{O}_{\Omega2} &= *\Omega_{\alpha\beta\mu}\phi^{;\mu} \mathring{\nabla}^{\alpha} \mathring{\nabla}^{\beta}\phi, \\ \mathcal{O}_{q1} &= q_{\alpha\beta\mu}\phi^{;\mu} \varphi^{;\alpha}\phi^{;\beta}\phi_{;\nu} \mathring{\nabla}^{\mu} \mathring{\nabla}^{\nu}\phi, \\ \mathcal{O}_{q3} &= q_{\alpha\beta\mu}\phi^{;\mu}\phi^{;\alpha}\phi^{;\beta}\phi^{;\rho}\phi^{;\sigma} \mathring{\nabla}_{\rho} \mathring{\nabla}_{\sigma}\phi. \end{aligned}$$

L₃ – systematic construction (HD)

• 5 Combinations that lead to 2nd order EL-eqs:

$$\begin{split} L_{3}^{(1)} &= \tilde{G}_{3}^{(1)}(\phi, X)(\tilde{\mathcal{O}}_{3} - \tilde{\mathcal{O}}_{1}) = 2\tilde{G}_{3}^{(1)}(\phi, X)\phi_{;\alpha}L^{\mu}{}_{\beta\mu} \left[\mathring{\nabla}^{\alpha}\mathring{\nabla}^{\beta}\phi - g^{\alpha\beta}\mathring{\Box}\phi \right] , \\ L_{3}^{(2)} &= \tilde{G}_{3}^{(2)}(\phi, X)(\tilde{\mathcal{O}}_{4} - \tilde{\mathcal{O}}_{6}) = \tilde{G}_{3}^{(2)}(\phi, X) \left(g_{\alpha\beta,\mu} - \Gamma^{\lambda}{}_{\mu\alpha}g_{\lambda\beta} - \Gamma^{\lambda}{}_{\mu\beta}g_{\alpha\lambda}\right) \left[\phi_{;\rho}g^{\mu\alpha}\mathring{\nabla}^{\rho}\mathring{\nabla}^{\beta}\phi - \phi^{;\mu}\mathring{\nabla}^{\alpha}\mathring{\nabla}^{\beta}\phi \right] , \\ L_{3}^{(3)} &= \tilde{G}_{3}^{(3)}(\phi, X)(\tilde{\mathcal{O}}_{5} - \tilde{\mathcal{O}}_{2}) = \tilde{G}_{3}^{(3)}(\phi, X)\phi^{;\mu} \left(g_{\mu\beta,\alpha} - \Gamma^{\lambda}{}_{\alpha\mu}g_{\lambda\beta} - \Gamma^{\lambda}{}_{\alpha\beta}g_{\mu\lambda}\right) \left[\mathring{\nabla}^{\alpha}\mathring{\nabla}^{\beta}\phi - g^{\alpha\beta}\mathring{\Box}\phi \right] , \\ L_{3}^{(4)} &= \tilde{G}_{3}^{(4)}(\phi, X)(\tilde{\mathcal{O}}_{10} - \tilde{\mathcal{O}}_{9}) = \tilde{G}_{3}^{(4)}(\phi, X)\phi^{;\alpha}\phi^{;\nu} \left(g_{\nu\rho,\mu} - \Gamma^{\lambda}{}_{\mu\nu}g_{\lambda\rho} - \Gamma^{\lambda}{}_{\mu\rho}g_{\nu\lambda}\right) \left[\phi^{;\mu}g^{\beta\rho} - \phi^{;\beta}g^{\rho\mu}\right]\mathring{\nabla}_{\alpha}\mathring{\nabla}_{\beta}\phi , \\ L_{3}^{(5)} &= \tilde{G}_{3}^{(5)}(\phi, X)(\tilde{\mathcal{O}}_{11} - \tilde{\mathcal{O}}_{7}) = \tilde{G}_{3}^{(5)}(\phi, X)\phi^{;\rho}\phi^{;\nu} \left(g_{\nu\rho,\mu} - \Gamma^{\lambda}{}_{\mu\nu}g_{\lambda\rho} - \Gamma^{\lambda}{}_{\mu\rho}g_{\nu\lambda}\right) \left[\phi^{;\alpha}g^{\mu\beta}\mathring{\nabla}_{\alpha}\mathring{\nabla}_{\beta}\phi - \phi^{;\mu}\mathring{\Box}\phi\right] , \end{split}$$

-(1)

Pairwise compensation

$$\tilde{\mathcal{O}}_{8} \text{ and } \tilde{\mathcal{O}}_{12} \text{ do not appear}$$
Can rewrite in basis of irreducible pieces
$$I = \begin{pmatrix} 1^{3} \\ 3^{2} \\ 4^{3}$$

L₃- systematic construction (HD)

• Total Lagrangian:

$$\begin{split} L_{\text{STKGB}} &= \mathring{R} + \mathring{L}_2 + \mathring{L}_3 + \sum_{a=1}^5 L_3^{(a)} \\ &= \mathring{R} + G_2(\phi, X) - G_3(\phi, X) \mathring{\Box} \phi + \mathcal{O}_1 G_3^{(1)}(\phi, X) + \mathcal{O}_2 G_3^{(2)}(\phi, X) + \mathcal{O}_3 G_3^{(3)}(\phi, X) \\ &+ \mathcal{O}_4 G_3^{(4)}(\phi, X) + \mathcal{O}_5 G_3^{(5)}(\phi, X) \,, \end{split}$$

- Recipe to go to higher order in Q, e.g. $N_Q = 2$ \implies tedious
- Also could go to higher i, e.g. i = 4 expect terms like $Q(\mathring{\nabla}\mathring{\nabla}\phi)^2$, $(\mathring{\nabla}\mathring{\nabla}\phi)\mathring{\nabla}Q$
- Potentially more general dependency on non-metricity-invariants in the generic functions.
 N_Q arbitrary

Cosmology: flat FLRW

• Consider:
$$L_{\text{ST-Horn}} = \sum_{i=4}^{5} \mathring{L}_i + G_{\text{ST}}(Q_i, \phi, X, I_i, J_i) + L_{\text{STKGB}}$$

• Ensure metric and connection respect the same symmetries: Killing vector fields Z_{ζ} where $\zeta = \{1, ..., m\}$

Solve
$$(\mathcal{L}_{Z_{\zeta}}g)_{\mu\nu} = 0$$
 $(\mathcal{L}_{Z_{\zeta}}\Gamma)^{\lambda}{}_{\mu\nu} = 0$
 $ds^{2} = -N(t)^{2}dt^{2} + a(t)^{2}(dr^{2} + r^{2}d\vartheta^{2} + r^{2}\sin^{2}\vartheta d\varphi^{2})$
 $g_{\mu\nu} = -n_{\mu}n_{\nu} + h_{\mu\nu}$ $n_{\mu} = (-N, 0, 0, 0)$
 $Q_{\rho\mu\nu} = 2F_{1}n_{\rho}n_{\mu}n_{\nu} + 2F_{2}n_{\rho}h_{\mu\nu} + 2F_{3}h_{\rho(\mu}n_{\nu)}$ $F_{i} = F_{i}(t)$ \Rightarrow 3 branches and $\Omega_{\lambda\mu\nu} = 0$

Conclusion

- Formulated scheme for potential extensions of Horndeski interactions in ST gravity
- Constructed the most general L₂
- Constructed an extension belonging to L_3 class up to $N_Q = 1$
- Future directions could include:
 - Check $c_T = 1$ constraints
 - Cosmological perturbations; strong coupling in f(Q) around one branch in cosmology
 - Scalarized BHs
 - Radiative stability

Thank you!